The four Pt-N distances found in the complex are not significantly different from one another. There is no evidence of any trans effect. The average Pt-N distance of 2.057 (6) Å is similar to those reported for *cis*and trans-dichlorodiamineplatinum(II)  $(2.05 \text{ (4)} \text{ Å})$ .<sup>12</sup> The nitrogen-platinum-nitrogen angles are within the usual range found in other ethylenediamine chelate compounds. Though there seem to be small differences in the nitrogen-carbon distances, the differences are not significant. A measure of the degree of nonplanarity of the ethylenediamine chelate ring is the dihedral angle of the two planes defined by the two carbon atoms and one of the donor nitrogen atoms, and the two carbon atoms and the other donor nitrogen in the ring. This dihedral angle is found to be 53" and *55'* in the two chelate rings of the complex. The analogous angles in  $[Co(C_2H_8N_2)_3]Br_3.0.5H_2O^{13}$  are 43-51<sup>°</sup>.

From the anomalous dispersion studies the absolute configuration of the complex is found to be A. This absolute configuration assignment is based on the following convention.<sup>14</sup> The planes each determined by the central metal atom and the two donor atoms of the bidentate ligand in a cis-bis-bidentate complex define a twofold axis and helix. A right-hand helix is labeled  $\Delta$  while a left-hand helix is labeled  $\Lambda$ . The conformations of the chelate rings are labeled according to a similar convention.<sup>14</sup> When the chelate ring is viewed along the axis through the two donor atoms the sequence of bonds in the chelate ring defines a static screw. If it is right-handed the conformation is labeled **6.** The mirror image is labeled A.

**(14)** *Y.* Saito, "Spectroscopy and Structure of Metal Chelate Compounds," K. Nakamoto and P. J. McCarthy, Ed., John Wiley & Sons, New .<br>York, N. Y., 1968, p 26.

The conformations of the two chelate rings, based on this convention, are  $\lambda \lambda$ . This is different from that of  $(+)$  -d-cis-dicyanobis (ethylenediamine) cobalt (III). where the absolute configuration of the chelate rings is opposite that of the  $C_2$  chirality of the complex ion.<sup>15</sup> However, it is analogous to the configuration found in  $tris(ethylenediamine) cobalt(III)$  chloride<sup>16</sup> and tris-(I-propylenediamine)cobalt(III) bromide. **l7** This combination of the absolute configuration of the complex and the conformations of the chelate rings is also the thermodynamically stable species in solution according to the conformational analysis of Corey and Bailar.<sup>18</sup> In investigations of tris(ethylenediamine)chromium(III) complexes Raymond, et  $al$ ,<sup>19</sup> found other combinations of absolute configuration of the complex and the conformations of the ethylenediamine chelate rings. They explained the existence of otherwise thermodynamically less favorable conformations on the basis of more extensive hydrogen bonding between the amine hydrogens and the cyanide groups. The fact that there are very few strong hydrogen bonds in the present structure is consistent with that explanation.

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(18) E. J. Corey and J. C. Bailar, Jr., *J. Ameu. Chem.* Soc., **81,** 2620 (1959). (19) K. N. Raymond, P. W. R. Corfield, and J. **A.** Ibers, *Iizoug. Chem., 7,*  842 (1968).

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## The Crystal Structure **of** Potassium **Trichloroammineplatinate(I1)** Monohydrate'"

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The crystal structure of potassium trichloroammineplatinate(II) monohydrate, K[PtCl<sub>3</sub>NH<sub>3</sub>]  $\cdot$ H<sub>2</sub>O, has been determined at room temperature by three-dimensional X-ray crystallographic techniques using intensity data measured by scintillation counter. After correction for absorption, refinement by the full-matrix least-squares method, using anisotropic temperature factors, converged to an *R* factor of 0.056. The compound crystallizes in an orthorhombic unit cell;  $a = 4.572 \pm 0.002$ ,  $b = 21.025 \pm 0.009$ ,  $c = 8.125 \pm 0.004$  Å with four molecules per unit cell and possible space groups Cmcm or Cmc2<sub>1</sub>. The latter is preferred. The anion is planar and the three Pt-Cl bond lengths of  $2.321 \pm 0.007$ ,  $2.315 \pm 0.027$ , and  $2.300 \pm$ 0.014 *b,* after correction for thermal motion, are equal within one standard deviation.

#### Introduction

Substitution reactions of platinum(I1) complexes with ammine and chloro ligands are often cited as examples of the so-called *trans* effect in metal complexes.<sup>2</sup> Thus the reaction of ammonia with the trichloroammineplatinate(II) ion gives only  $cis$ -dichlorodiammineplatinum(II), the chloro ligand opposite another chloro ligand is replaced in preference to that opposite ammonia, and chloride is said to have a higher trans effect than ammonia.

An accurate determination of the structure of the  $[PLC1_3NH_3]$ <sup>-</sup> ion should reveal whether the apparent

<sup>(13)</sup> K. Nakatsu, *B~LII. Chem.* Soc. *Jap.,* **36,** 832 (1962).

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*<sup>(2)</sup>* F. Basolo and R. G. Pearson, *Progv. Inorg. Chem.,* **4,** 381 (1963).

chemical difference between Pt-Cl bonds *cis* and *trans*  to an ammine group is reflected in a difference in the corresponding Pt-Cl bond lengths.

**A** single crystal electron diffraction structure determination of anhydrous potassium trichloroammineplatinate(I1) found the Pt-C1 bond *trans* to chloride to be 0.03 A longer than the one *trans* to ammonia; however, this difference has an estimated standard deviation of  $\pm 0.03$  and so cannot be regarded as significant. The present determination was undertaken at the suggestion of the late Professor R. E. Rundle to try to determine the Pt-C1 bond lengths in the [PtCls-  $NH<sub>3</sub>$  - anion with greater accuracy.

## Experimental Section

Crystals of  $K[PtCl_3NH_3] \cdot H_2O$  which grow in large flat plates were generously supplied by Dr. D. S. Martin. Attempts were made to grind spherical crystals using an apparatus similar to that described by Bond.4 The crystals were rather fragile and often shattered as they approached a suitable size; the best crystal finally selected was ellipsoidal in shape,  $0.08 \times 0.08 \times 0.16$  mm.

Preliminary Laue and precession photographs showed the cell to be orthorhombic, with  $a = 4.572 \pm 0.002$ ,  $b = 21.025 \pm 0.009$ ,  $c = 8.125 \pm 0.004$  Å, obtained by micrometering precession photographs of a very small crystal to  $\pm 0.04$  mm. (Errors reflect internal consistency; the camera was previously calibrated using a NaCl crystal,  $a = 5.6399$  Å, for Mo K $\alpha$ ,  $\lambda$  0.7107 Å). This is in agreement with Bokii, *et al.*,<sup>3</sup> who found  $a = 20.6$ ,  $b = 8.0$ , and  $c = 4.5$  Å, and with Lyashenko,<sup>5</sup> who reported axial ratios  $1.2910:1.0:0.8404$  (corresponds to our  $a:2b:3c = 1.294$ :  $1.0:0.8626$ ).

The density, measured by flotation in a mixture of methylene iodide and 1,2-dibromoethane, was found to be  $3.181 \pm 0.003$ g/ml from four separate determinations. Calculated values of 3.04 for **4** molecules of KPtC13NH3, and 3.19 for 4 molecules of  $KPtCl<sub>3</sub>NH<sub>3</sub>·H<sub>2</sub>O$ , indicated that the crystal examined had the latter formula.

Absences of  $h + k = 2n$  for *hkl* and  $l = 2n$  for *hOl* reflections gave the possible space groups as Cmcm (no. 63,  $D_{2h}^{17}$ ), Cmc $2_I$ (No. 36,  $C_{2v}^{12}$ ), or C2cm (No. 40,  $C_{2v}^{16}$ ) (Ama2 in standard setting).

**A** complete octant of three-dimensional intensity data out to  $\sin \theta / \lambda = 0.75$  was collected for the crystal at room temperature using molybdenum radiation  $(\lambda = 0.7107 \text{ Å})$ . A General Electric XRD-5 X-ray unit with single crystal orienter and scintillation counter was used with the moving-crystal-moving-counter technique and  $(\theta, 2\theta)$  coupling, a 3<sup>°</sup> take-off angle, and a 2<sup>°</sup> diffracted beam aperture fitted with a zirconium foil filter. Single crystal orienter settings were precalculated using a program sco-6 written for the IBM 7074 computer by D. E. Williams. A 200 sec scan at  $1^{\circ}/\mathrm{min}$   $(2\theta)$  was used for each reflection, followed by a repeat scan for background with angle *w* offset by 1.8' for about 40% of reflections. During the period of 8 days required for these measurements, three standard reflections (0, 0, 16; 0, 4, 6; 2, 1, 2) were checked daily and the intensities of these reflections remained constant over a period of 8 days to within  $\pm 3\sigma$ , which was  $\pm 3.5\%$  of total counts. A total of 830 reflections were measured.

**A** graph of measured background counts *us. 0* was plotted which showed the background to vary smoothly with  $\theta$ ; consequently, the background correction for each reflection was obtained by interpolation from a table as a function of **0.** In view of the large absorption coefficient  $(\mu = 203 \text{ cm}^{-1})$ , an absorption correction was applied approximating the crystal by a 26-face truncated prism and using a 216-point grid for the numerical

**(3) G. B.** Bokii, B. K. Vainshtein, and **A. A.** Babareko, *Izo. Akad. Nauk SSSR, Old. Khim. Nauk,* **6, 667 (1951).** 

**(4) W. L.** Bond, *Rev. Sci. Instrum.,* **22, 344 (1951).** 

*(5)* **W. N.** Lyashenko, Tr. *Inst. Kristallogv., Akad. Nauk SSSR,* **9, 335 (1954).** 

integration.B **A** correction for the presence of noncharacteristic radiation along radial reciprocal lattice lines was made by a similar method to that described by Williams and Rundle,7 adapted by J. E. Benson of this Laboratory. The intensity *vs.* **e**  profiles of the  $[0\ 0\ 1]^*$  and  $[1\ 0\ 3]^*$  lines were used.

In the later stages of least-squares refinement, a weighting

factor w was applied to each reflection, calculated from  
\n
$$
\sigma F = 1/\sqrt{w} = \left[\sqrt{(I + \sigma I)} - \sqrt{I}\right]/(Lp\sqrt{A})
$$
\n
$$
(\sigma I)^2 = \left[C_{\rm T} + C_{\rm B} + (K_{\rm T}C_{\rm T})^2 + (K_{\rm B}C_{\rm B})^2 + (K_{\rm S}C_{\rm S})^2(K_{\rm A}C_{\rm P})^2\right]/A^2
$$

where  $C_T$ ,  $C_B$ , and  $C_S$  are the counts for total peak, background, and streak, and  $C_P = C_T - C_B - C_S$ ;  $K_T$ ,  $K_B$ ,  $K_S$ , and  $K_A$  are nonstatistical proportional errors in total peak, background, streak counts, and absorption correction and were taken to be 0.04, 0.04, 0.1, and 0.1, respectively; *A* is the absorption correction (actually a transmission factor);  $L$  is the Lorentz factor; and  $\dot{p}$  is the polarization factor. The terms  $\sigma I$  and  $\sigma F$  represent the estimated errors in intensity and structure factor, respectively.

Atomic scattering factors were taken from Cromer and Waber,8 and in the later stages of refinement were modified for the effects of anomalous dispersion using the coefficients of Cromer **.Q** 

## Structure Determination

Since there are only four molecules per unit cell, and no systematically weak reflections were observed, the platinum atoms must lie on special positions  $c$  of Cmcm  $(0, y, \frac{1}{4})$ , *a* of Cmc2<sub>1</sub>  $(0, y, z)$ , or *b* of C2cm  $(x, y, \frac{1}{4})$ , all of which have only one variable parameter  $y$  (in the last two space groups a second parameter may be arbitrarily chosen for 1 atom). For platinum only structure factor calculations were carried out with values of y incrementing by 0.025 from  $y = 0$  to  $y =$ 0.25. The *R* factor showed a pronounced minimum at  $y = 0.15$  and least-squares refinement improved this coordinate to 0.157. For convenience this was altered to the equivalent position of  $y = 0.343$ . In Cmcm the platinum has site symmetry mm, and coordinates for the two nonrelated chlorine atoms Cl(1) (on both mirror planes) and  $Cl(2)$  (one of the mirror related *trans* pair) were readily derived.

At this point structure factor calculations gave an  $R$  factor of 0.18. Attempts to place the potassium ion in the available fourfold sites of Cmcm were unsuccessful. **A** three dimensional observed Fourier synthesis calculated with Cmcm symmetry, and based on the platinum and chlorine positions, revealed the nitrogen atom at the expected position on the mirror plane at *z*   $=$   $\frac{1}{4}$ , and also two mirror-related pairs of peaks at  $x = 0$  whose peak heights corresponded to half-potassium and half-oxygen atoms. Two interpretations of the half-height peaks were considered. The four potassium and four oxygen atoms could occupy the eightfold  $sites f (0, y, z)$  in Cmcm, or  $c (x, y, z)$  in C2cm, in a random manner while the  $[PtNH<sub>3</sub>Cl<sub>3</sub>]$ <sup>-</sup> ion retains full mm symmetry. An alternative interpretation is that the true space group is  $Cmc2<sub>1</sub>$ , which does not require a mirror plane at  $z = \frac{1}{4}$  and consequently full K and O atoms may be placed one on either side of the now nonexistent mirror plane. In this case the  $[PtCl<sub>3</sub>NH<sub>3</sub>]$ <sup>-</sup> ion is only required to be planar.

(6) W. R. Busing and **H.** Levy, *Acto Cvystallogu.,* **10, 180 (1957).** 

- **(7) D. E.** Williams **and** R. **E.** Rundle, *J. Amer. Chrm. Soc.,* **86, 1660 (1964).**
- **(8) D. T.** Cromer and J. T. Waber, *Acta Crrstallogv.,* **18, 104 (1965).**
- **(9)** D. **T.** Cromer, *ibid.,* **18, 17 (1965).**

TABLE I REFLECTIONS REMOVED FROM FINAL REFINEMENT *h k I Fo Fo h k I F, Fo*  0 2 *0* 125 74 0 2 1 265 303 0 8 0 122 *87* **1** 1 **6** 15 **94** 

0 14 0 68 33 3 9 *7* 91 39

Structure factors calculated from the platinum, chlorine, and nitrogen atom positions, and with potassium and oxygen positions with a population parameter of 0.5, in space group Cmcm gave an agreement factor  $R = 0.15$ , which rapidly converged to 0.13 with isotropic temperature factors, and to 0.068 (0.058 omitting zero reflections) with anisotropic temperature factors for all atoms. At this stage, the  $Cl(2)$  atom anisotropic thermal vibration component perpendicular to the molecular plane was very large, suggesting that the

sides of the hypothetic mirror plane at  $z = \frac{1}{4}$ , at first neglecting the imaginary part of the platinum scattering factor. A slight improvement in  $\Sigma w \Delta^2$  was obtained together with small shifts of low significance, in atom positional parameters.

Since  $Cmc2<sub>1</sub>$  is a polar space group, the presence of anomalously scattering atoms could lead to systematic errors in atomic position,<sup>10</sup> particularly as only one octant of data was measured. Duplicate refinements were carried out, first assuming the octant was  $+h$ ,  $+k$ , *+l* and then  $-h$ ,  $-k$ ,  $-l$ . Ten cycles of refinement showed different behavior. The  $+h$ ,  $+k$ ,  $+l$  case converged and on the last cycle the largest changes were for the  $U_{11}$  components of atoms  $Cl(2)$  and  $Cl(3)$ , which the U<sub>11</sub> components of atoms C1(2) and C1(3), which amounted to  $\frac{1}{30}$  of one standard deviation. The  $-h$ , amounted to  $\frac{1}{s_0}$  of one standard deviation. The  $-h$ ,<br>-  $k$ ,  $-l$  case was consistently divergent and could not be

TABLE **I1** 

MULTIPLIED BY 10. RUNNING INDEX IS  $k$ , FOLLOWED BY  $F_{\text{obsd}}$  AND  $|F_{\text{calcd}}|$ 

OBSERVED AND CALCULATED IN THE RESIDENT CONTINUES IN THE RESIDENT OF THE RESID  $\begin{array}{cccc} 17 & 75 & 83 \\ 19 & 197 & 193 \\ 21 & 33 & 496 \\ 23 & 31 & 864 \\ 25 & 87 & 84 \\ 27 & 4 & 222 \\ 29 & 84 & 832 \\ 21 & 35 & 222 \\ 31 & 35 & 222 \\ 31 & 21 & 288 \\ 5 & 159 & 167 \\ 7 & 197 & 190 \\ \end{array}$ *3 243 2+3 5 \*5* 41 **7 163 15L**  とうにしゃっしゃ こうかんか こうみつか こった アルカイ しゅんしょう かくりょう こうこうしょう うんしゅう しゅうしょう あいしょう こうしんのうてき こうしん こうりん こうしょう しょうし しゅんき こうしょ こうしょう しょうしょう しょうしゅう こうしょう しょうしゅうしゅ しょうしゅう こうしょう の4136457 しゅんのシアア ウェスクアール ちりろうち うつうし しきこん ラス りょうしょう ライトラ きょうしょう しょうしょう こうしょう こうしゅう こうしゅう こうしゅう こうしょう こうしょう こうしょう こうしょう こうしょう しょうしょう しょうしゅ ひとん こしいしょうかん こうかん こうこうかん こうしょう こうしゅう こうこう こうしょう こうしょう こうしょう こうこうしょう こうこう こうしょう こうこうしゅう こうかん こうかん こうかん こうこうかん いちょうこう 3824424114288828930  $\frac{1}{3}$ <br>  $\frac{5}{7}$ <br>  $\frac{9}{11}$ <br>  $\frac{11}{11}$  $3882300$ **9** *2n7* **212**  li **30 Jb**  13 **169** *11"*  **15 Y5** *92*  **17 57** *I,* I? **133 133** *21* **27 L8 21 Ll 62 25 7) 19 <sup>21</sup>***b* **0 29 LO 5,**  . **H.** I **Li** *<sup>5</sup>* 1 **128 125 3 56** *55 5 b,* I1 *7* **I17** 113 *9 IC1* **LOP**  11 **18 133 13 12** 11 **15 \$3 iB**  22998410785512019200281124623559315 **9 166 ir3**   $\begin{array}{cccccccc} 11.2 & 30 & 24.9 & 10.1 & 11.2 & 10.$ 19 0 1 3 4 5 6 7 7 8 9 7 11 in1 101  $\begin{array}{l} 12122814\\ 2212329\\ 23195\\ 242195\\ 252195\\ 252195\\ 252195\\ 252195\\ 252195\\ 252111\\ 25211\\ 25211\\ 25211\\ 25211\\ 25211\\ 25211\\ 25211\\ 25211\\ 25211\\ 25211\\ 25222\\ 25211\\ 25221\\ 25221\\ 25211\\ 25221\\ 25211\\ 25221\\ 2521$ 1614102435389488829017235978317775435999317775 6 5 3 1 2 7 7 7 2 6 5 1 3 0 3 3 4 6 1 9 2 8 5 8 **27 81 80 2**<br>29 19 14 **\***<br>\* H \* 1 1 \* \* \* \* *<sup>29</sup>***19 lk**  1 **172 lbl** 

 $[PtCl<sub>3</sub>NH<sub>3</sub>]=$  ion was not strictly planar. Three cycles of refinement in space group C2cm, after arbitrarily displacing the  $Cl(2)$  atom out of the plane by  $0.3 \text{ Å}$ , did not improve the  $R$  factor, nor effectively cause any lowering of temperature factors.

If the water molecule and potassium ion occupy disordered sites, there is a possibility that the number of water molecules per potassium ion is less than one, especially in view of the similarity of the structure of this hydrated complex with the anhydrous compound. Therefore the population parameter for oxygen was made variable, but least-squares refinement led to a value corresponding to a full molecule of water per potassium ion, within the limit of error.

Refinement was then continued in space group Cmc-21, with full potassium and oxygen atoms on opposite refined to a stable set of parameters. This difference of behavior suggests that the measured octant was  $+h$ ,  $+k$ ,  $+l$ . However, a search through the calculated sets of structure factors revealed that for only 24 reflections was the difference  $|F_{hkl}| - |F_{hkl}|$  greater than 1.5 $\sigma$  for the measured value.

A comparison of the final  $+h$ ,  $+k$ ,  $+l$  set of parameters with those obtained in  $Cmc2<sub>1</sub>$ , neglecting  $\Delta f''$ for platinum, revealed a shift of  $+0.037$  (7) Å along the *z* direction for the potassium atom, relative to platinum. This shift is approximately that predicted by Cruickshank and McDonald.<sup>10</sup> Similar small shifts were observed for the other atoms, but in all other cases these were of the order of one standard deviation.

**<sup>(10)</sup>** D. **W.** J. **Cruickshank and W.** S. McDonald, **Acta** *Cvystallogr.,* **18, 9**  (1967).

The final *R* factor was 0.069 for all reflections, and 0.056 for nonzero reflections. Six reflections (Table I) were removed from the final refinement as they appeared to contain large random errors  $(\sqrt{w}\Delta > 3)$ . A threedimensional difference Fourier synthesis revealed no significant features which could be attributed to hydrogen atoms. The observed and calculated structurefactor amplitudes are given in Table 11. Table 111

TABLE **I11**  FRACTIONAL ATOMIC COORDINATES FOR REFINEMENT IN Cmc2<sub>1</sub> AND Cmcm<sup>a</sup>

		-Cmc2--		Cmcm∙		
	$\boldsymbol{\gamma}$	z	y	z		
Pt.	0.34338(3)	$1/\epsilon$	0.34342(3)	1/4		
Cl(1)	0,4536(2)	0.2506(28)	0.4531(2)	1/4		
C1(2)	0.3373(9)	0.5311(39)	0.3402(2)	0,5316(6)		
Cl(3)	0.3425(6)	$-0.0306(21)$	(0.3402(2))	$-0.0316(6)$		
N	0.2467(9)	0.2277(63)	0.2449(9)	1/4		
$\mathbf{K}$	0,0395(3)	0.4640(9)	0.0400(3)	0.4595(9)		
$\circ$	0.0756(11)	0.1374(31)	0.0738(9)	0.1300(29)		
	<sup>a</sup> Standard deviation on last figure in parentheses.			All atoms		

are at  $x = 0$ .



Figure 1.--Molecular geometry in the  $[PtCl_3NH_3]$ <sup>-</sup> anion. Standard deviation of last figure in parentheses.

after correction for anisotropic thermal motion by the method of Busing and Levy.<sup>11</sup>

**A** clear choice between the results obtained from space groups Cmcm and  $Cmc2<sub>1</sub>$  is not easily made. Refinement in both space groups leads to a similarly low  $R$  factor, and in the case of  $Cmc2<sub>1</sub>$  the presence of a polar dispersion effect does not completely clarify the situation. The Hamilton *R* factor ratio test suggests that the Cmc2<sub>1</sub> refinement is highly significant  $( \theta_1 =$  $0.68/0.63 = 1.080$  using the weighted *R* factor for all

TABLE IV ANISOTROPIC THERMAL PARAMETERS (x **io4 ia)** FOR REFINEMENT

	ANISOTROPIC LHERMAL PARAMETERS (XIU' A') FOR REFINEMENT		
	IN Cmc <sub>2</sub> <sup>1</sup> AND Cmcm, WITH STANDARD DEVIATIONS IN PARENTHESES <sup>2</sup>		



*a* The form of the temperature factor is  $\exp[-2\pi^2(U_{11}h^2a^{*2} + U_{22}k^2b^{*2} + U_{33}l^2c^{*2} + U_{23}klb^{*}c^{*})].$ 

TABLE V

	MOLECULAR GEOMETRY OF $[PtCl3NH3]$ - ANION					
	-Cmc21 refinement-		Cmcm refinement-			
Atoms	Dist, Å	Dist. A. cor for thermal vibration	Dist. A	Dist. Å. cor for thermal vibration		
$Pt-C1(1)$	$2.317 \pm 0.007$	$2.321 \pm 0.007$	$2.306 \pm 0.005$	$2.310 \pm 0.005$		
$Pt-Cl(2)$	$2.288 \pm 0.027$	$2.315 \pm 0.027$	$2.289 \pm 0.006$	$2.315 \pm 0.006$		
$Pt-C1(3)$	$2.280 \pm 0.014$	$2.300 \pm 0.014$				
$Pt-N$	$2.042 \pm 0.019$	$2.060 \pm 0.019$	$2.07 \pm 0.02$	$2.09 \pm 0.02$		
Atoms $Cl(1)-Pt-Cl(2)$ $Cl(1) - Pt - Cl(3)$		Angle, deg	Angle, deg			
		$93.1 \pm 0.6$ $91.7 \pm 0.1$ $90.6 \pm 0.5$ A				
	$Cl(2)-Pt-N$	$91.9 \pm 0.9$				
				$88.3 \pm 0.1$		
	$Cl(3)-Pt-N$	$84.4 \pm 0.7$				

lists the final atom position parameters, and Table IV gives the final atom vibration parameters obtained from the refinements in space groups Cmcm and Cm $c2<sub>1</sub>$ .

## **Discussion**

Figure 1 shows the intramolecular geometry of the  $[PtNH<sub>3</sub>Cl<sub>3</sub>]<sup>-</sup>$  ion, which is crystallographically required to be planar in  $Cmc2<sub>1</sub>$ . Table V gives bond lengths and angles for the parameters obtained from the Cm $c2<sub>1</sub>$  refinement, and for the Cmcm refinement where  $\mathrm{K}^{+}$  and the oxygen from a water molecule are assumed to be in disordered positions. Distances are also given reflections,  $\mathcal{R}_{11,782,005} = 1.035$ . However, the reliability of this test under similar circumstances of pseudo symmetry has been questioned.<sup>12</sup>

However, it can be concluded from the structure that Cmcm cannot be the true space group symmetry of an individual unit cell since this would lead to the improbable situation of having two potassium ions per platinum in some cells, and two water molecules per platinum in others. Impossible contact distances arise from this situation. The crystal could, of course, be disordered in such a way that the two possible  $Cmc2<sub>1</sub>$  ar-

**<sup>(11)</sup> W. R. Busing and H. A. Levy, Ada C~yslallogr., 17, 142 (1964)** 

**<sup>(12)</sup> A. Whitaker and J. W. Jeffrey,** *ibid.,* **P8, 984 (1967).** 



Figure 2.-Crystal packing of K[PtCl<sub>a</sub>NH<sub>3</sub>] *H*<sub>2</sub>O, projected down the *a* axis. Open circles show atoms at  $x = 0$ , shaded circles are atoms at  $x = \frac{1}{2}$ . Distances shown are in angström units.

rangements, one with  $K^+$  on the left and  $H_2O$  on the right of the pseudo mirror, the other *vice versa,* both occur equally in the crystal. This would lead to an apparent crystal symmetry of Cmcm. Refinement of such a structure would lead to a set of atom positions which would be the average of the two  $Cmc2<sub>1</sub>$  arrangements and would naturally have Cmcm symmetry. These positions would not necessarily correspond to the actual geometry of any individual molecule, however. It should be possible to distinguish between a Cmcm (disordered) lattice and an ordered  $Cmc2<sub>1</sub>$  lattice using the Pt anomalous dispersion effect. However, although some indication was gained from the parallel refinements of  $F_{hki}$  and  $F_{hki}$  data that the former rather than the latter had been measured, the magnitude of the effect is not sufficiently large for an unequivocal assignment to be made.

Nevertheless, we favor the results of the  $Cmc2<sub>1</sub>$ refinement for both crystallographic and chemical reasons. The Cmc2<sub>1</sub> refinement did successfully converge to the structure presented here from a number of different starting models. Also, the resulting crystal lattice bears a striking resemblance to that of  $K[PLCl_3 \cdot$  $C_2H_4$ .  $H_2O$ ,  $13$  whose structure has recently been reinvestigated. The formation of this ordered lattice (discus;ed later) may provide additional stability over a disor lered structure.

In view of the expected interactions between the  $[PtCl<sub>3</sub>NH<sub>3</sub>]$  anion and the potassium ions and water molecules, it was not felt that the  $PtCl<sub>3</sub>NH<sub>3</sub>$  group would behave as an independent rigid body in the lattice for a molecular vibration analysis to be applicable. In the  $Cmc2<sub>1</sub>$  refinement the very high correlation between corresponding parameters of  $Cl(2)$  and  $Cl(3)$ unfort inately lowers the accuracy of the two *trans*  Pt-Cl bond lengths. The smaller errors for the Cmcm refinement cannot be taken at face value, either, for the reasons given above. The anomalously high component of the thermal vibration normal to the molecular plane for these two atoms also increases the uncertainty in these Pt-C1 bond lengths. It is apparent that there is no significant difference between the Pt-C1 bond lengths with or without any correction for thermal motion, although the limits of error are unavoidably high.

(13) M. Black, R. H. B. hlais, and P. G. **Owston,** private communication.

This result appears to conflict with the prediction of the polarization theory of GrinbergI4 for the *trans*  effect. This theory assumes that a strongly *trans*directing group actually weakens the metal-ligand bond opposite to itself by an electrostatic polarization mechanism, thus enabling *trans* substitution to occur more easily than in the *cis* position. In the  $[PtCl_3NH_3]^$ complex ion, the Pt-C1 bond *trans* to chlorine is expected to be weaker, and presumably longer, than the Pt-C1 bond *trans* to ammonia. If there is any weakening, it is apparently not accompanied by any appreciable lengthening (>0.05 **A)** of the Pt-C1 bond.

An alternative explanation of the *trans* effect by Chatt<sup>15</sup> and Orgel<sup>16</sup> which suggests  $\pi$ -bond stabilization of a five-coordinate transition state does not seem applicable in this case because of the very low  $\pi$ bonding ability of the ammine and chloro ligands.

**A** similar result has been recently found by Milburn and Truter" for the *cis* and *trans* isomers of dichlorodiammineplatinum(II), where the Pt-Cl bonds are not significantly different in the two compounds:  $2.32 \pm 0.01$  Å for *trans*,  $2.33 \pm 0.01$  Å for *cis*.<br>In contrast, an appreciably longer Pt–Cl bond (2.422)  $\pm$  0.01 Å for *trans*, 2.33  $\pm$  0.01 Å for *cis.* 

 $\pm$  0.009 Å) *trans* to a hydrido ligand has been observed in **trans-hydridochlorobis(diphenylethy1phosphine)**  platinum(II).<sup>18</sup> The hydrido ligand shows a very much greater *trans* effect than chloride or ammonia,<sup>19</sup> and the lack of any significant difference in Pt-C1 bond lengths in amminechloroplatinum(I1) complexes may show that the effect is too small to be detected. It is interesting to note that there is no significant difference in the rate of exchange of isotopically labeled chloride ions with the *cis* and *trans* chlorine atoms in the [PtC13-  $NH<sub>3</sub>$ ]-ion.<sup>20</sup>

In **cis-dichlorobis(triethylphosphine)platinum(II),**  long Pt–Cl bond lengths of  $2.364$  and  $2.388$  Å have been observed.21 However, no lengthening of the Pt-C1 bond occurs in a corresponding *trans* compound, *trans*dichlorobis **(trimethy1phosphine)platinum** (I I), **22** Furthermore, the bond length of  $2.294 \text{ Å}$  found is very close to the average length, 2.289 A, of the two uncorrected Pt-Cl(2) and Pt-Cl(3) distances in the present structure; this suggests there is no  $cis$  effect.<sup>23</sup>

The crystal packing and intermolecular contacts are shown in Figure 2 and Table VI. The crystal structure of the hydrated complex is quite similar to that of the anhydrous compound determined by Bokii, *et* al.3 In the latter structure, the potassium is on the center of symmetry at the origin in space group Cmcm *(a* =

- (17) G. **H.** W. Milburn and M. R. Truter, *J. Chem. Sac., A,* 1609 (1966).
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- *J. Chem.* Soc., 2207 (1961). (20) T. *S.* Elleman, J. **W.** Reishus, and D. *S.* Martin, *J. Amev. Chem. Sac.,*
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**(23) P.** G. **Owston,** J. M. Partridge, and J. M. Rowe, *Acla Cvystailogv.,* **lS, 246** (1960).

<sup>(14)</sup> **A.** A. Grinberg, *Acta Physicochim. URSS,* **8,** 573 (1935).

<sup>(15)</sup> J. Chatt, L. **A.** Duncanson, and L. M. Venanzi, *J. Chem.* Soc., **<sup>4456</sup>** (1955).

<sup>(16)</sup> L. E. Orgel, *J. Inovp. Nucl. Chem.,* **2,** 137 (1956).





*a* Atom at 0,  $-y$ ,  $\frac{1}{2} + z$ . *b* Atom at  $\frac{1}{2}$ ,  $\frac{1}{2} + y$ , *z. <i>c* Atom at  $\frac{1}{2}$ ,  $\frac{1}{2} - y$ ,  $\frac{1}{2} + z$ . **4** Atom at  $\frac{1}{2}$ ,  $\frac{1}{2} - y$ ,  $z - \frac{1}{2}$ . **4** Atom at 0,  $y, 1 + z$ .

4.19,  $b = 17.6$ ,  $c = 8.84$  Å), and is surrounded by eight chlorine atoms in a distorted cubic coordination. In the monohydrate, the potassium is displaced from the origin to a site which has six chlorine atoms at the corners of a trigonal prism and two water molecules in the centers of two of the rectangular faces. Thus water can be relatively easily lost from the lattice without much disruption of the structure, particularly in the Cmc21 model where the water molecules all lie above each other along the *x* axis, but the featureless difference Fourier synthesis and normal thermal parameters for oxygen confirm that the crystal examined here was fully hydrated. The potassium-oxygen distances (2.76 (2) and 2.80 (2) Å) are normal, the three potassium-chlorine distances **(3.26** *(2))* **3.37** (l), **3.39** (1) A) are somewhat longer than is commonly found  $(3.20-3.30 \text{ Å})$ .<sup>23</sup>

The water molecule has an approximately tetrahedral environment, consisting of two potassium ions and two chlorine atoms [(angle Cl-O-Cl  $97^\circ$ , K-O-K  $104^\circ$ )]. The oxygen-chlorine distance of **3.05** (2) **A** is slightly less  $(0.15 \text{ Å})$  than the sum of the van der Waals radii for oxygen and chlorine, and it is tempting to suggest that the water hydrogen atoms lie on these directions with a weak hydrogen bond between them. This is also a favorable electrostatic configuration of the oxygen atom with respect to the potassium ions.

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# **The Crystal and Molecular Structure of**  Bis(3,3',5,5'-tetramethyldipyrromethenato)nickel(II)<sup>1</sup>

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The crystal and molecular structure of bis(3,3',5,5'-tetramethyldipyrromethenato)nickel(II), Ni(C<sub>13</sub>N<sub>2H16</sub>)<sub>2</sub>, has been determined by single-crystal X-ray diffraction methods. The crystal data are: space group A2/a  $(C_{2h}^6$ , no. 15); cell dimensions  $a = 22.352 + 0.007 \text{ Å}$   $b = 7.337 + 0.005 \text{ Å}$   $c = 31.496 + 0.008 \text{ Å}$   $s = 114.70 + 0.05^{\circ}$   $V$  $a = 22.352 \pm 0.007$  Å,  $b = 7.337 \pm 0.005$  Å,  $c = 31.496 \pm 0.008$  Å,  $\beta = 114.70 \pm 0.05^{\circ}$ ,  $V = 4692.7$  Å<sup>3</sup>; density 1.29  $\pm$  $0.02$  g/cm<sup>3</sup> (by flotation), 1.294 g/cm<sup>3</sup> (calcd for  $Z = 8$ ). The intensities of 1616 reflections collected on a manually operated counter diffractometer with Cu *Ka* radiation were used in solution and refinement. The structure was solved by conventional Patterson and Fourier methods and refined by full-matrix least-squares methods to final unweighted and weighted residuals of 0.068 and 0.097, respectively. The molecule possesses distorted  $D_2$  symmetry, and the dihedral angle between the least-squares planes formed by the two ligands (including the nickel atom in each case) is 76.3". The nickel-nitrogen bond lengths (1.948 (7)-1.957 (8) **8)** are identical within experimental error. The bond lengths among the ring atoms indicate that electron delocalization occurs over each planar ring system.

#### **Introduction**

Metal complexes of dipyrromethenes are of interest because of their close relation to the biologically important metalloporphyrins. Several papers have appeared in the literature in which the preparation, spectra, magnetic properties, and possible structures of metal complexes of substituted dipyrromethenes are discussed. Much of this work centers on ligands containing methyl substituents in the 5 position of each ring. Porters dealt with bivalent metal complexes *(Co,* Ni, Cu, Zn, Cd, Pd) and concluded on the basis of models that a planar configuration of the ligands was impossible due to steric repulsions between the *5-* and 5'-methyl groups

on one ligand and those on the other. Even with hy drogen atoms in the *5,5'* positions, there would probably be sufficient repulsion that true planarity  $(D_{2h})$  could not be achieved.

Recent work by Ferguson and Ramsey4 and Ferguson and West<sup>5</sup> favors a tetrahedral or distorted tetrahedral arrangement of the ligands. However, their work as well as that of Murakami and Sakata<sup>6</sup> seems to indicate that removal of one or both substituents in the *5* positions on the rings allows a twisting toward planarity, especially in a coordinating solvent.

In view of the uncertainty as to the detailed structure of this type of complex and as a basis for our further

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**<sup>(4)</sup>** J. E. **Ferguson and** C. **A. Ramsey,** *ibid.,* **5222 (1965).** 

**<sup>(5)</sup>** J. **E. Ferguson and B.** *0.* **West,** *ibid., A,* **1565 (1966).** 

**<sup>(6)</sup>** *Y.* **Murakami and K. Sakata,** *1x0~~.* **Chin.** *Acta,* **2, 273 (1968).**